

TWO CLASSES OF MAGNETOTAIL DIPOLARIZATION FRONTS OBSERVED BY MAGNETOSPHERIC MULTISCALE MISSION A STATISTICAL OVERVIEW



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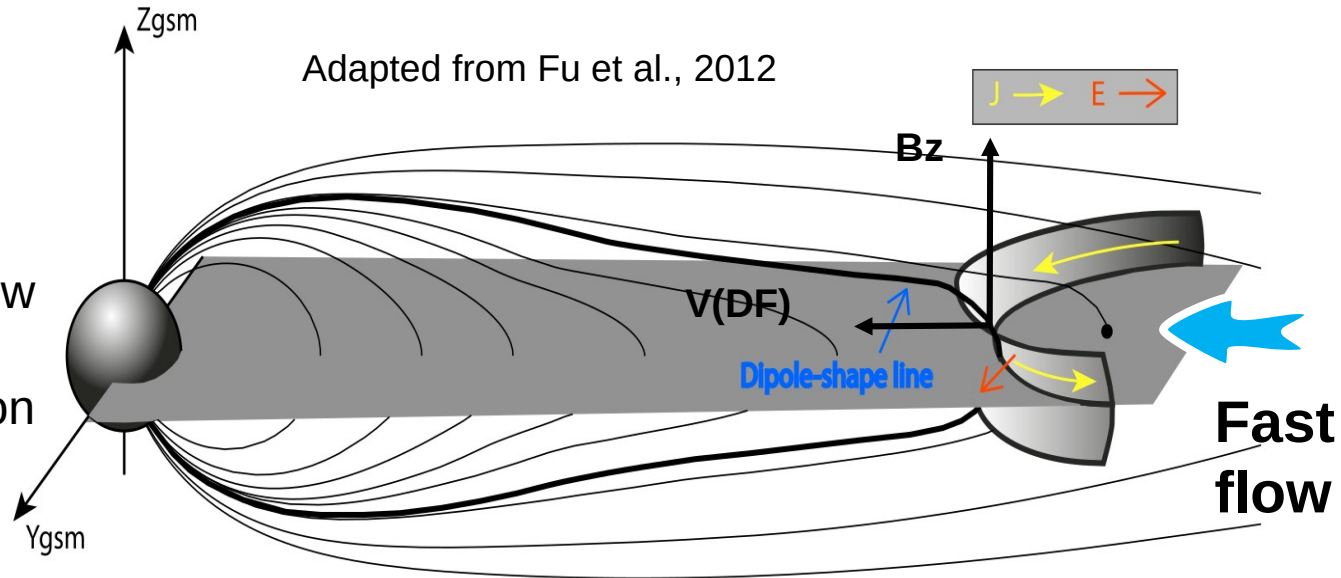
S. W. Alqeeq, O. Le Contel, P. Canu, A. Retinò, T. Chust, L. Mirioni, A. Chuvatin, R. Nakamura, N. Ahmadi, F. D. Wilder, D. J. Gershman, Yu. V. Khotyaintsev, P.-A. Lindqvist, R. E. Ergun, J. L. Burch, R. B. Torbert, S. A. Fuselier, C. T. Russell, H. Y. Wei, R. J. Strangeway, K. R. Bromund, D. Fischer, B. L. Giles, and Y. Saito



Dipolarization Front (DF)



- Sharp increase in the northward component of the magnetic field (Zdirection)
- associated with a fast plasma flow
- can be generated by reconnection or kinetic ballooning interchange instability.
- DF corresponds to a boundary between a relatively cold and dense plasma at rest and a hot tenuous fastly moving plasma.



Typical DF properties

$V(DF) \sim 200$ km/s

Thickness ~ 500 km

Crossing time ~ 2.5 s

[e.g., Runov et al., 2009 using THEMIS; Fu et al., 2012 using Cluster]

One MMS DF example

16:46:30-16:49:00 UT



Alqeeq et al., POP 2022

DF/fast flow properties

[e. g., Runov et al., 2009, Sergeev et al., 2009]

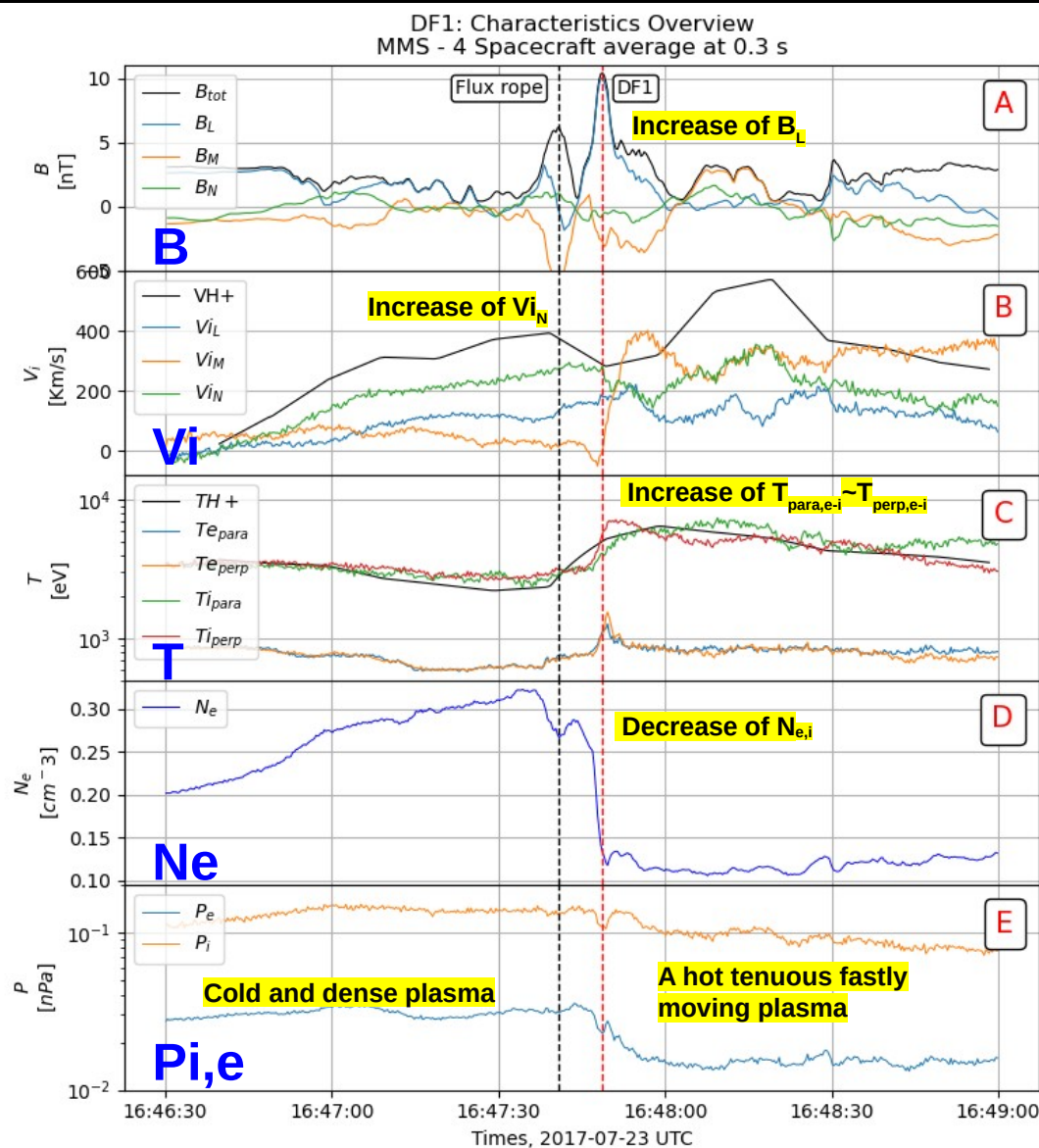
MVAB analysis on 4 s/c averaged data
between (16:47:45/16:48:00)

LMN frame of DF:

$L = (0.14, 0.63, 0.76)$

$M = (0.13, -0.78, 0.62)$

$N = (0.98, 0.01, -0.19)$



The main objective



- In order to extend these case study results, I have carried out a statistical analysis over the full 2017 magnetotail season.
- In particular, the energy conversion process and its homogeneity at electron scales are investigated

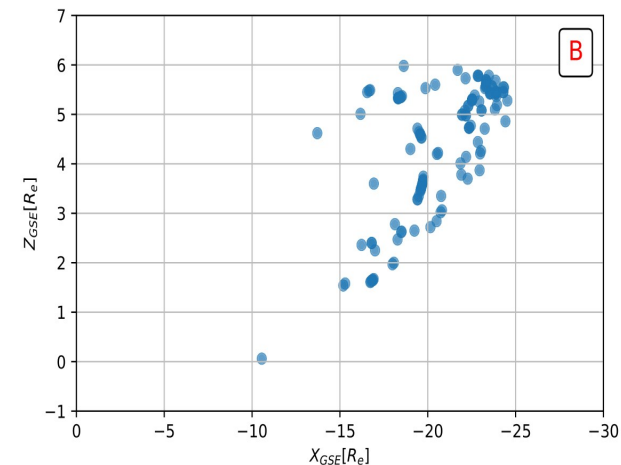
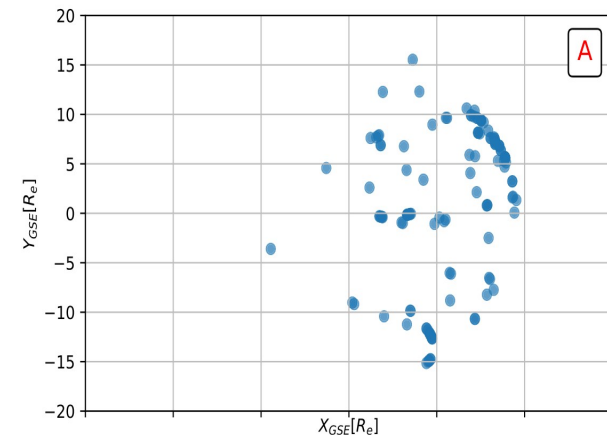
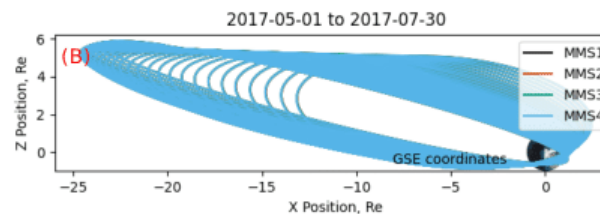
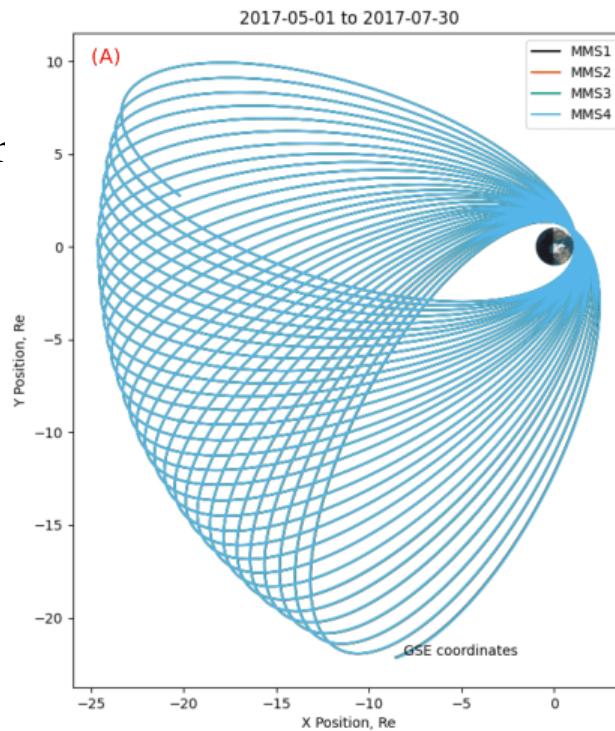


Selection criteria : A statistical study of DFs



Alqeeq et al., JGR 2023

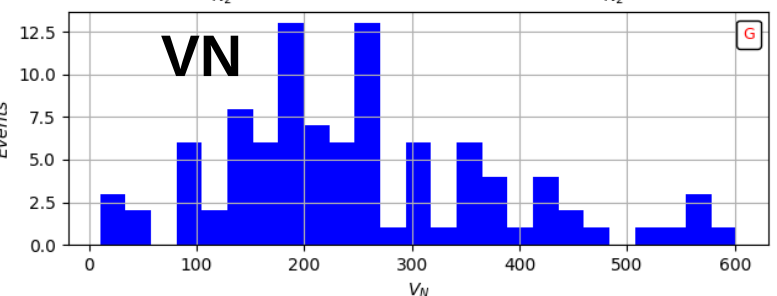
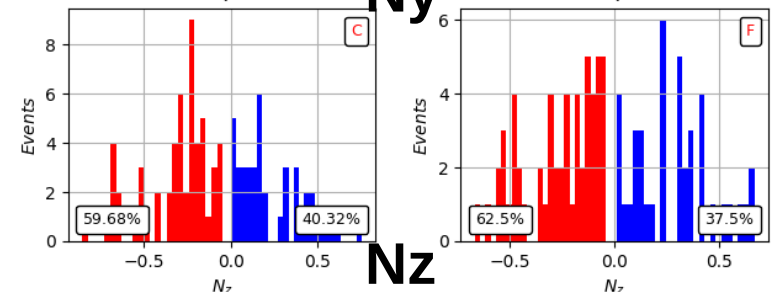
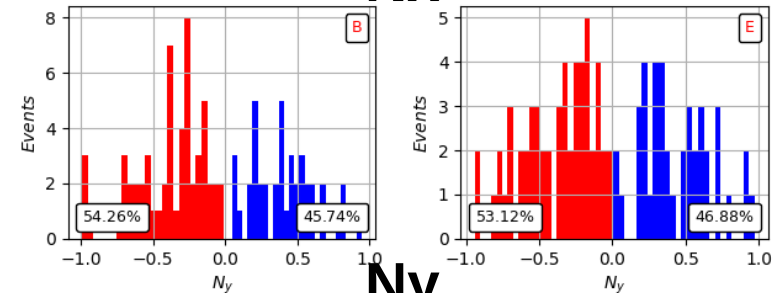
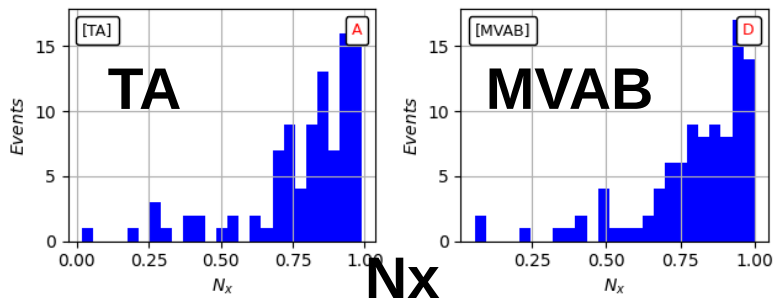
- More than 800 “possible DF” events detected near the Earth’s magnetotail equator ($|\mathbf{B}_x| < 5\text{nT}$), using an AIDApY tool requesting \mathbf{B}_z and \mathbf{V}_i increases and \mathbf{N}_e decrease.
- **This first automatic selection is then adjusted manually with the following criteria leading to only 132 DF events:**
- Burst mode (partmoms) data are available at least 30s before and after the DF. The head of the DF denotes the time t_0 .
- \mathbf{B}_z increase $> 5\text{ nT}$
- $\mathbf{V}_i > 150\text{ km/s}$
- $\mathbf{N}_{e,i}$ decrease
- $\mathbf{T}_{\text{para},e-i} \sim \mathbf{T}_{\text{perp},e-i}$ increases.



Methods



- MVAB and TA methods are applied on magnetic field data. They give similar front normal.
- MVAB was set to be:
 - L : always oriented northward
 - M: directed dawnward
 - N: earthward
- **Almost no preferential directions of propagation along y and z.**
- **N_y : ~ (+) 54% , ~ (-) 46%**
- **N_z : ~ (+) 60% , ~ (-) 40%**
- **V_N : ~200 km/s**

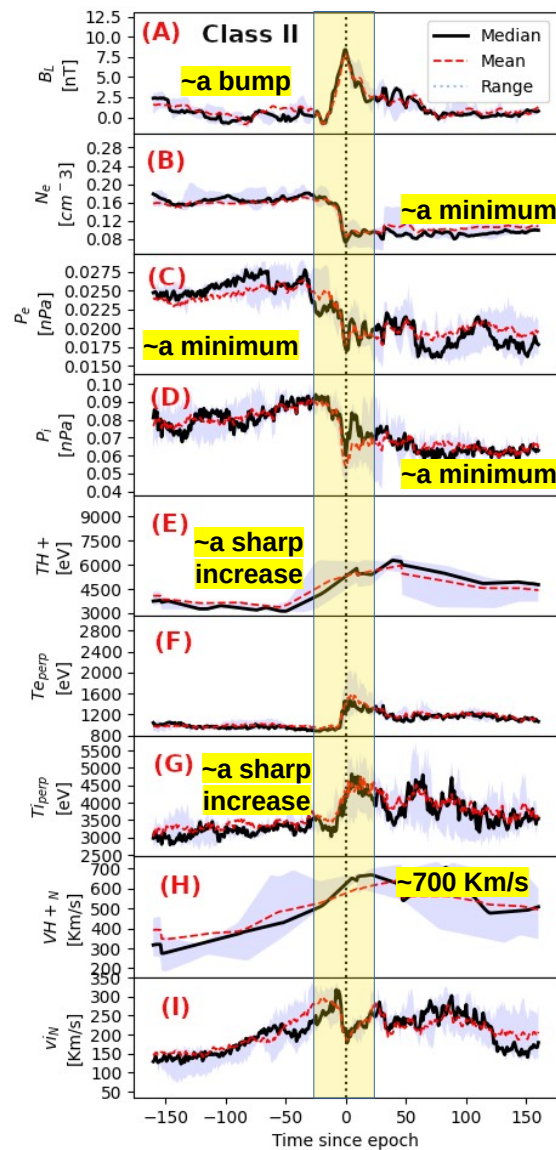
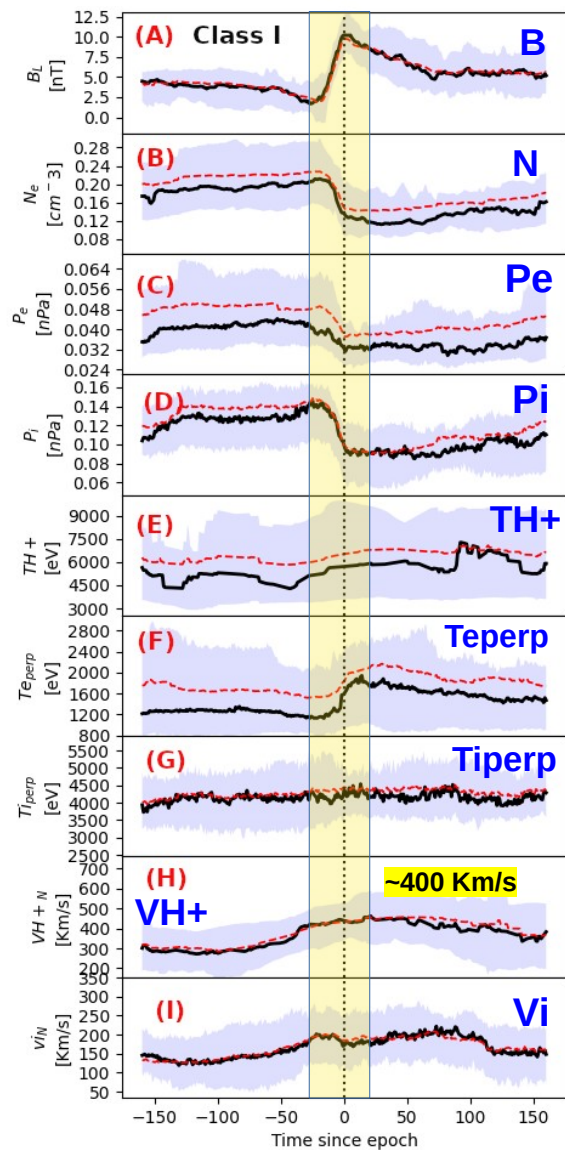


Two classes of magnetotail DFs



DFs

- Class I “classic type” (74.4%)** corresponds to a slow decrease of the magnetic field after the DF and is associated with smaller ion velocity and hotter plasma [e. g., Schmid et al.,2015; Huang et al.,2015; Yao et al.,2015; Zhong et al.,2019].
- Class II “new type” (25.6%)** has the same time scale for the rising and the falling of the magnetic field (a bump) associated with a minimum of ion and electron pressures and faster velocity [S. W ALQEEQ et al.,2022].

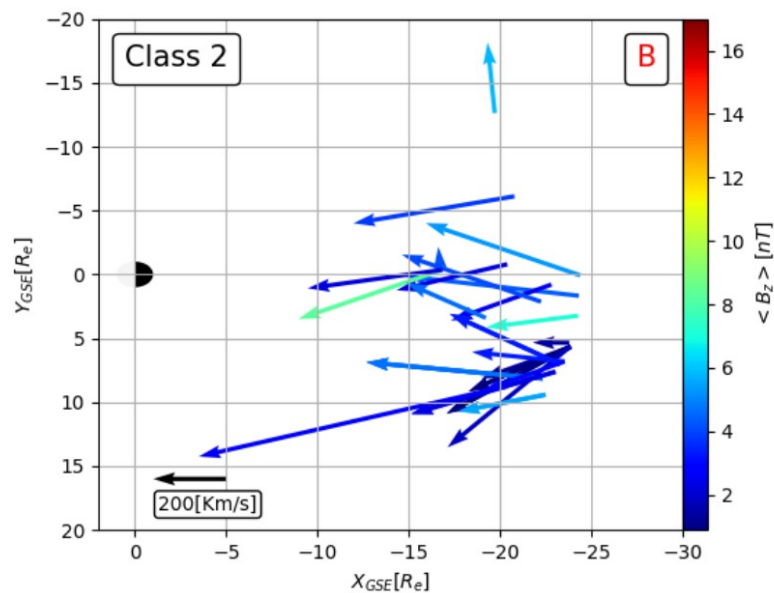
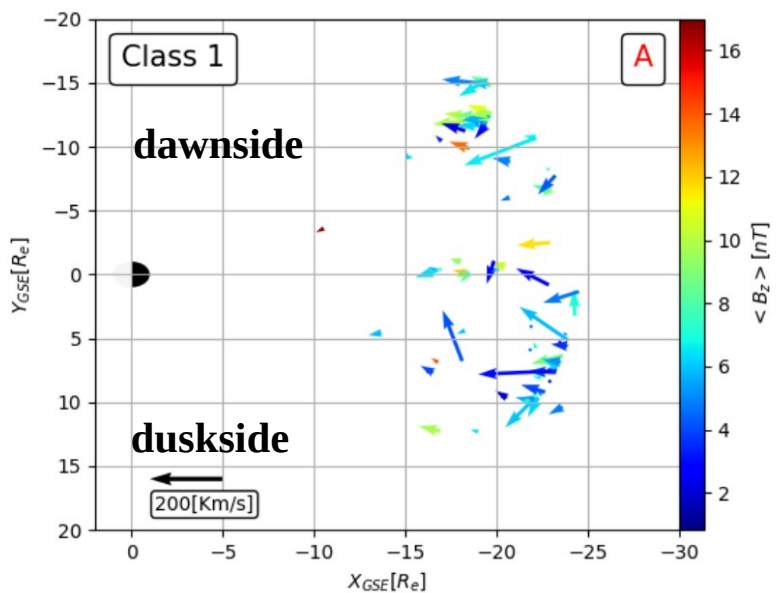


An overview of the Class I and Class II events



- **Class I** DF locations and propagations are relatively **random**
- **Class II** DF have preferentially **duskward locations** and propagations with **larger velocities**.

Equatorial XY GSE Plane



Current density comparisons

MMS - 4 Spacecraft average at 0.3s

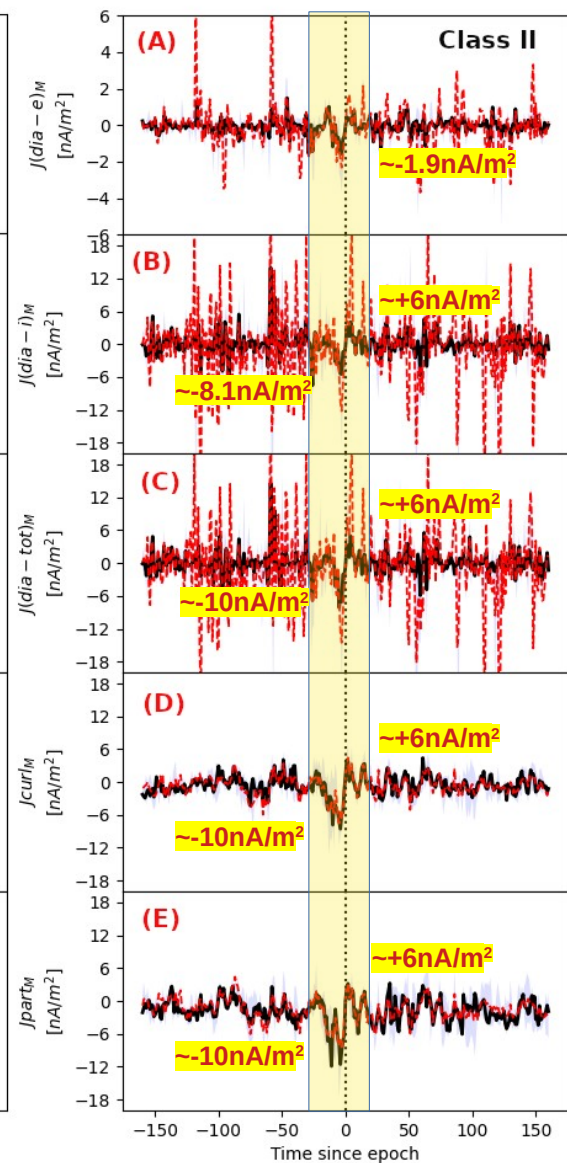
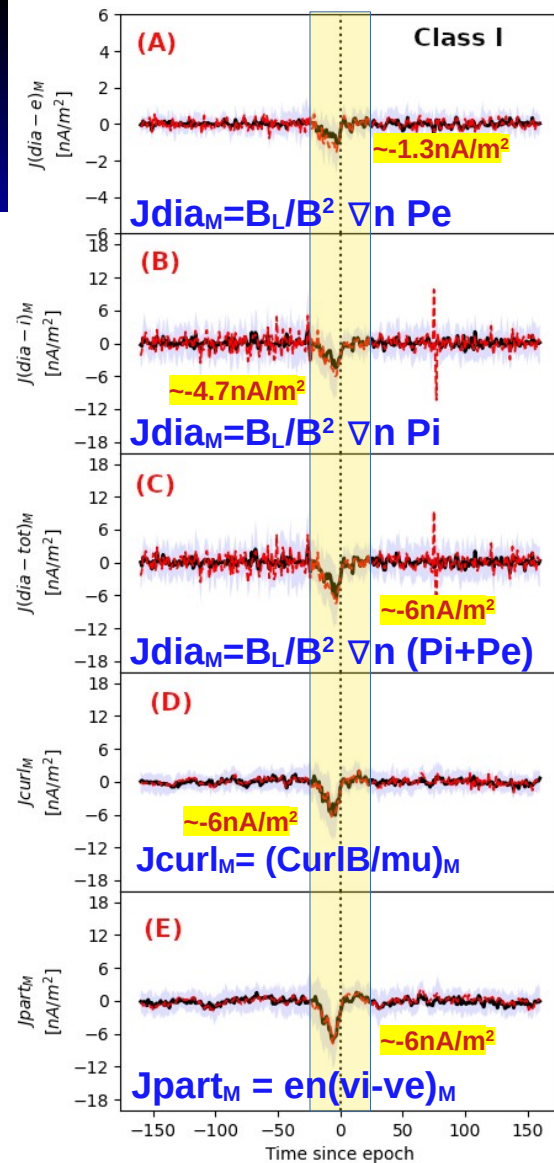


For both classes:

Small values but good agreement within $<10\text{nA/m}^2$

Ion diamagnetic current is dominant ($\sim 72\%$)

In Class II the reversal in J_{part_M} (E) & J_{curl_M} (D) is due to the reversal of the diamagnetic current (C), dominated by ions (B)



Ion Ohm's Law



$$\mathbf{E} + \mathbf{v}_i \times \mathbf{B}$$

Ideal ion
frozen-in

$$\frac{\mathbf{J} \times \mathbf{B}}{en}$$

$$-$$

Hall
electric
field

$$\frac{1}{en} \nabla \cdot \mathbf{P}_e$$

$$-$$

Electron
pressure
gradient

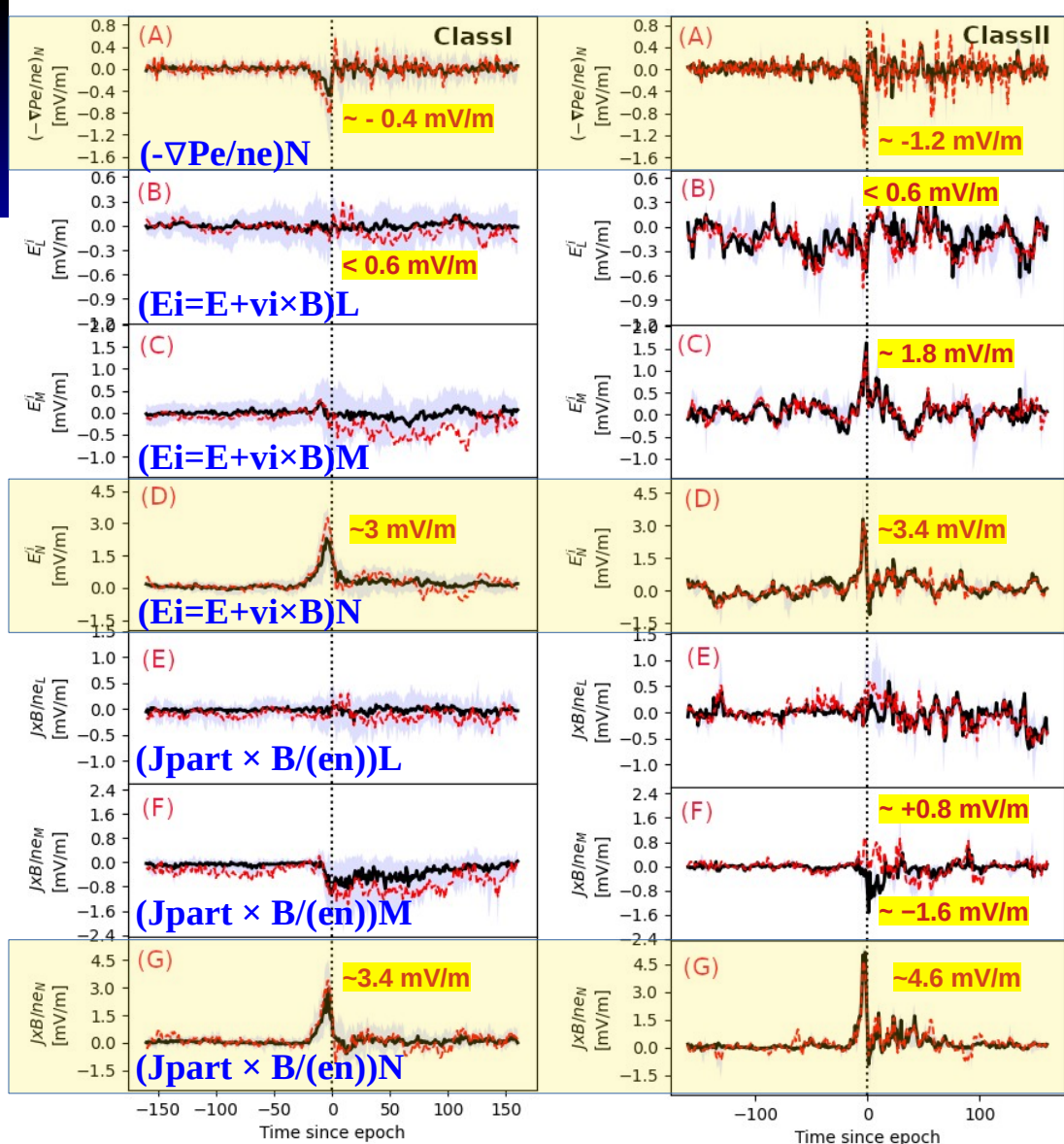
$$-\frac{m_e}{e} \frac{dv_e}{dt}$$

$$-$$

~~Initial
term~~

For both classes **in the N direction**, ions are **decoupled (D)** mostly by the **Hall electric field (G)** but electron pressure could also contribute (assuming non-zero curl) (A).

- For class II in the M direction, ions are decoupled (C), but the Hall field (median ~ -1.6 mV/m, mean $+0.8$ mV/m) suggests that the contribution from the electron pressure gradient could be quite large too ($\sim +3.4$ mV/m, $+1$ mV/m, respectively).

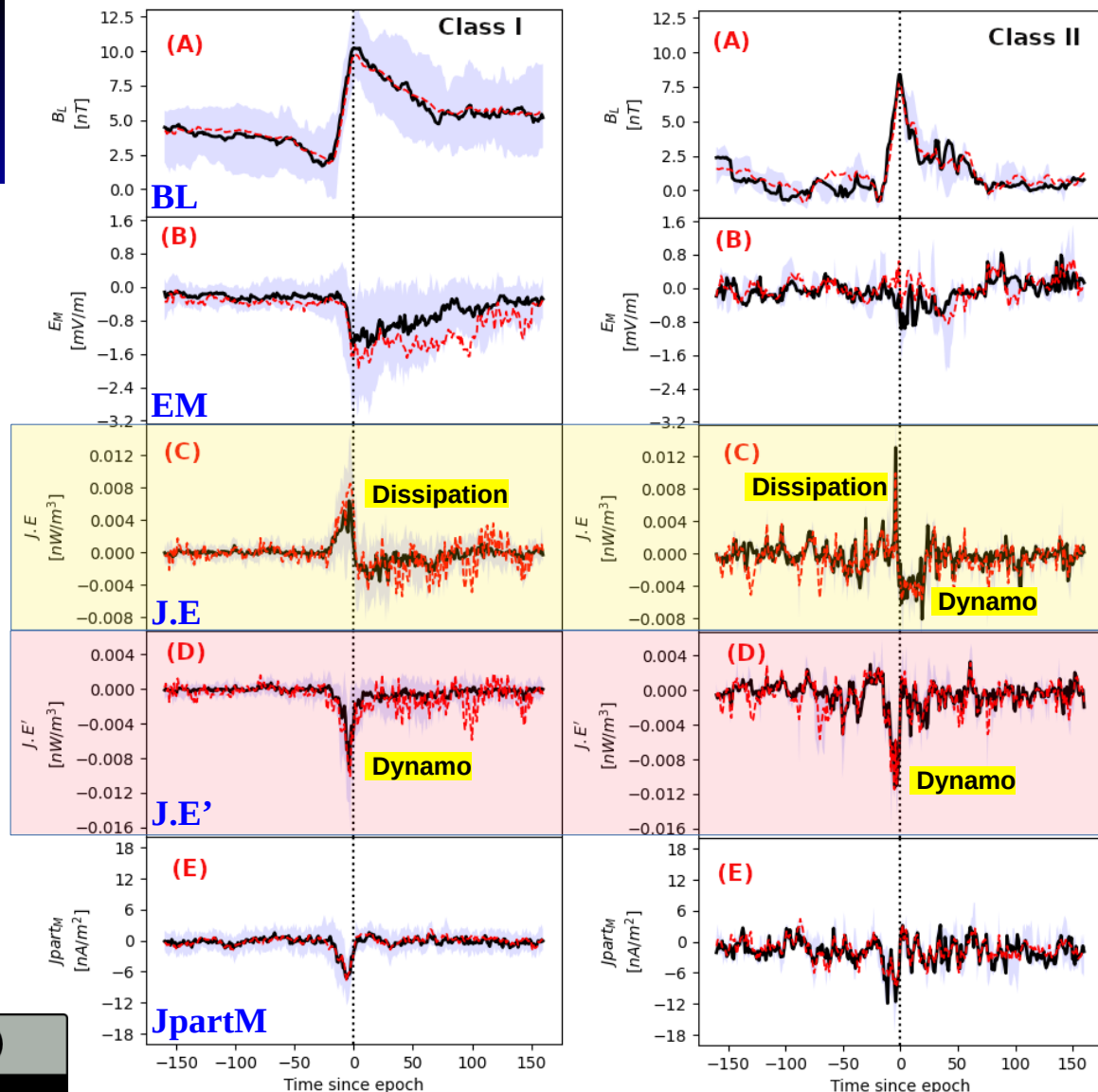


Energy conversion

MMS - 4 Spacecraft average at 0.3s



- For both classes ahead of DF, in the s/c frame, $\mathbf{J}_{part.E} > 0$ **Dissipation** (C) (energy is dissipated from the electromagnetic field to the particles).
- For class II behind DF, in the s/c frame, $\mathbf{J}_{part.E} < 0$ **Dynamo** (C) (the energy is transferred from the particles to the electromagnetic field).
- For both classes ahead of DF, in the ion & electron frames, $\mathbf{J}_{part.E'} < 0$ **Dynamo** (D) (energy goes from particles to field).
- In Class II the reversal of $\mathbf{J}_{part.E}$ (C), in the s/c frame is due to the reversal of the diamagnetic current.



Standard Deviation analysis for E' & Jpart



Compute the standard deviation (SD) normalized by its error bar:

$$SD(X)/\Delta X = \sqrt{\sum_{i=1}^4 (X_i - \langle X \rangle)^2 / 4} / \Delta X$$

$\langle X \rangle$: The four spacecraft average of the X component

ΔX : Respective estimated error bar

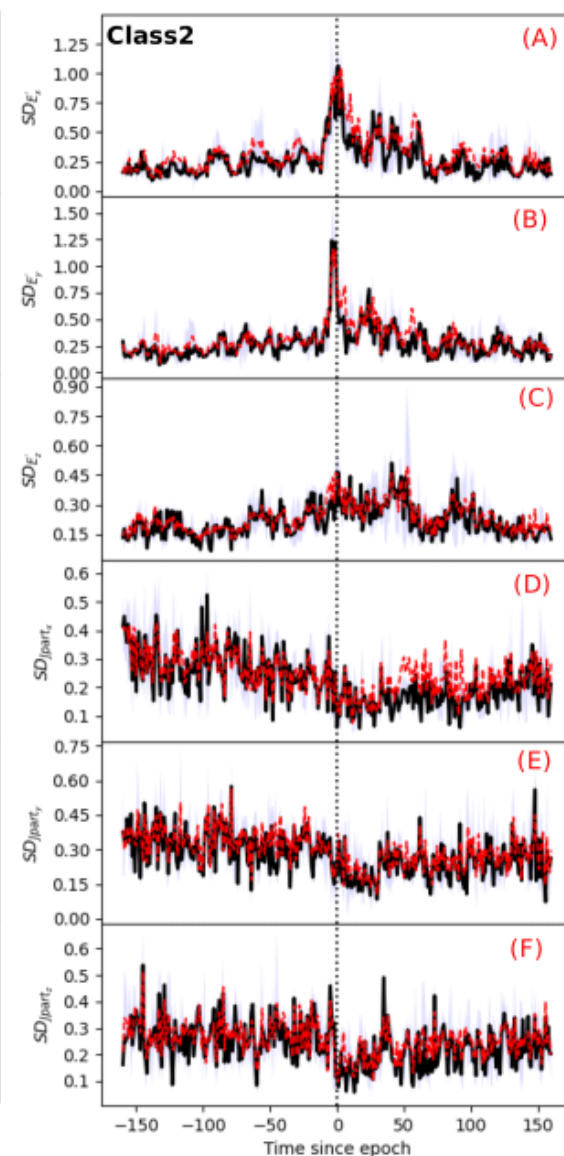
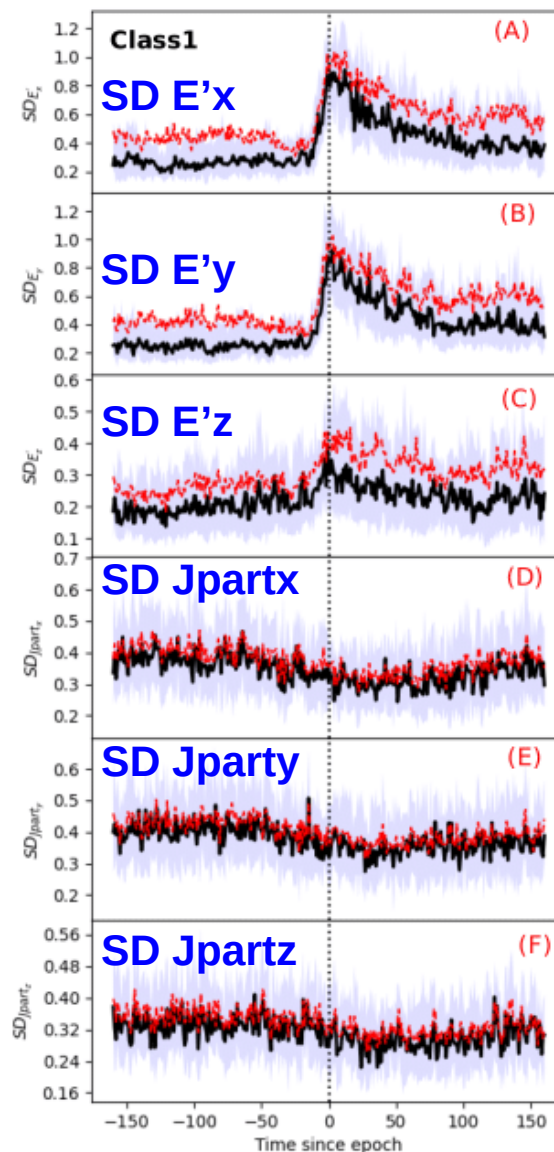
$\Delta E'$ ~1.7 mV/m

ΔJ ~6.8 nA/m²

- For both classes **normalized** SD of **E'** fields is about 1 for x and y components (A & B) whereas SD of current densities is always smaller than 1 for all components.

Thus:

- These statistical results confirm that **the non homogeneity comes from the E' field** as shown by Alqeeq et al. 2022 for six DF events.



Conclusion (I)



For the full magnetotail season of 2017 (132 DF events):

- **Class I “classic type”** (74.4%) corresponds to a **slow decrease of B** after the DF and is associated with smaller ion velocity and hotter plasma.
- **Class II “new type”** (25.6%) has a **bump B profile** associated with a minimum of ion and electron pressures and faster velocity as shown in Alqeeq et al. 2022, and it is found mostly on the duskside.
- For both categories we found a good agreement between current densities calculated from particles, Curl B and J_{diaM} (single S/C method).
- For both categories we found that **ions are mostly decoupled from the magnetic field by the Hall fields.**

The non- zero curl of the **electron pressure gradient term is also contributing to the ion decoupling and responsible for an electron decoupling** at DF.

Conclusion (II)



Summary of class I and class II signatures

Physical quantity	Class I	Class II
B_z or B_L	fast increase then slow decrease	fast increase then fast decrease "bump"
$P_{i,e}$ & $N_{i,e}$	monotonous decrease	minimum or "hole"
$\mathbf{J} \cdot \mathbf{E}$ (s/c frame)	>0 dissipation	>0 dissipation then <0 dynamo
$\mathbf{J} \cdot \mathbf{E}'$ (fluid frame)	<0 dynamo	<0 dynamo
Geometry	2D $E'_N \sim E'_M \sim 0.8 $ mV/m	2D $E'_N \sim E'_M \sim 1.2 $ mV/m
Homogeneity at electron scales $SD(E')$	No (>1)	No (>1)

The electron pressure gradient term could also contribute (assuming a non-zero curl) to the ion decoupling and lead to the electron decoupling.



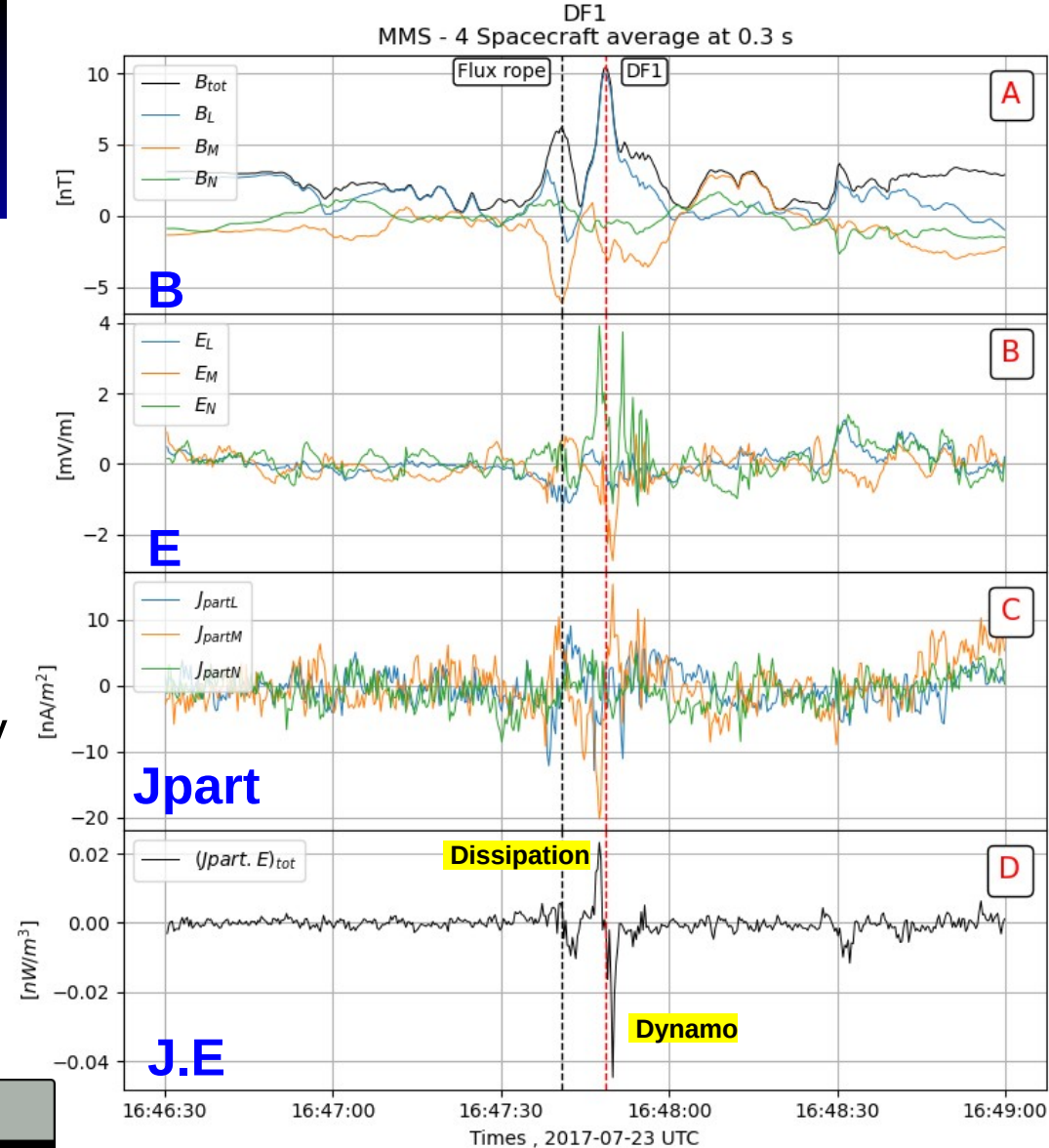
**Merci de votre
attention**

Energy conversion (I)



In s/c frame:

- “Dissipation” ahead of the front, the energy is transferred from the electromagnetic field to the particles
 $\mathbf{J} \cdot \mathbf{E} \sim +0.023 \text{ nW/m}^3$.
- “Dynamo” behind the front, the energy is transferred from the particles to the electromagnetic field
 $\mathbf{J} \cdot \mathbf{E} \sim -0.043 \text{ nW/m}^3$.
- Convective field ($\mathbf{E} < 0 \sim \mathbf{v} \times \mathbf{B}$) dominant
- $\text{Sign}(\mathbf{J} \cdot \mathbf{E}) \sim \text{sign}(J_{\text{part},M})$



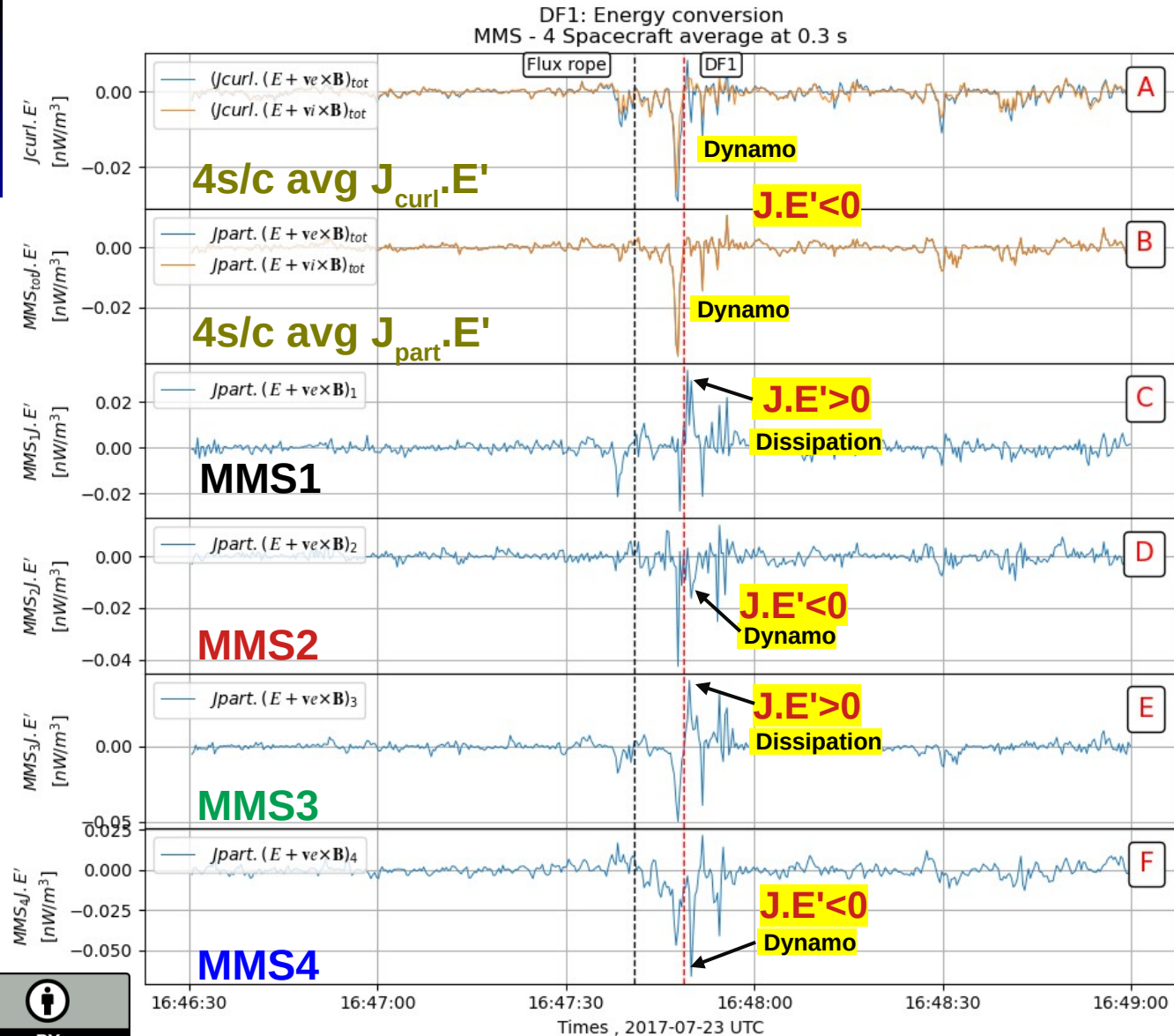
Energy conversion (II)

16:47:30-16:48:40 UT

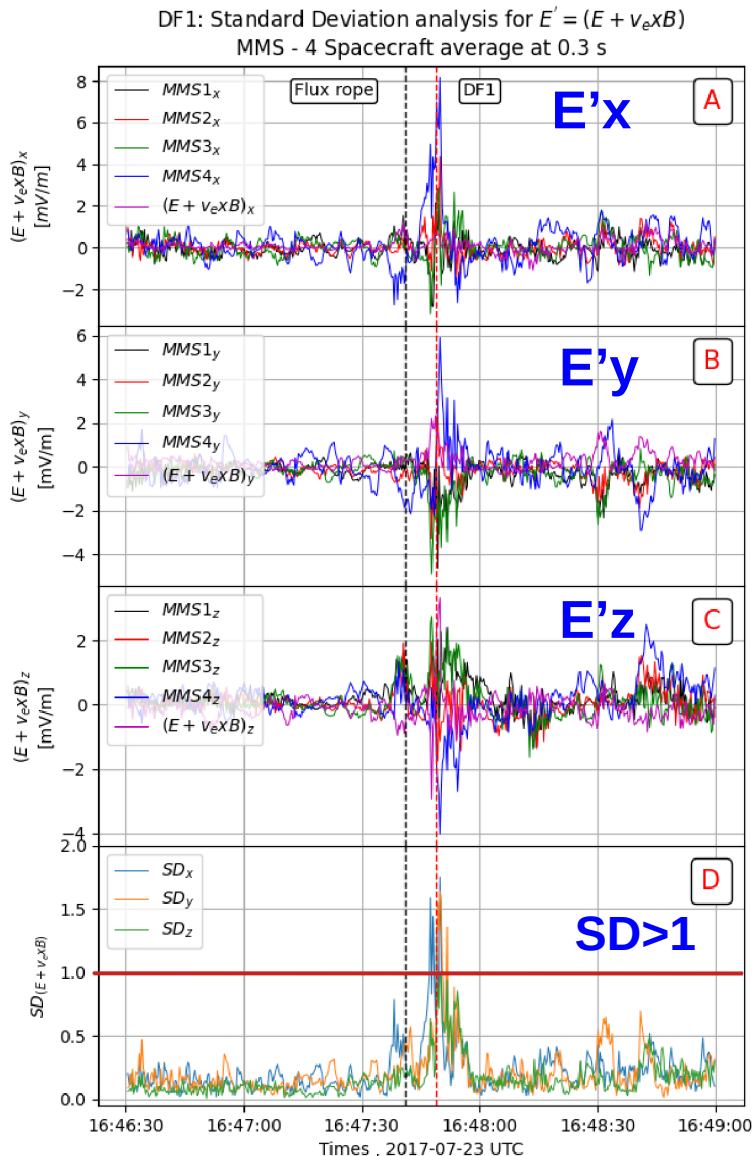


In Ion & electron frames:

- ⇒ Good confidence with all $\mathbf{J} \cdot \mathbf{E}' = \mathbf{J} \cdot (\mathbf{E} + \mathbf{v}_i \times \mathbf{B})$ calculations.
- $\mathbf{J} \cdot \mathbf{E}' > 0$, **Dissipation** (energy goes from field to particles) ~ after the DF (from single s/c MSS1, 3)
- $\mathbf{J} \cdot \mathbf{E}' < 0$, **Dynamo** (energy goes from particles to field) ~ at DF (from 4 s/c and all singles s/c)
- These results are consistent with [Yao et al., 2017].



Standard Deviation analysis for E' & Jpart

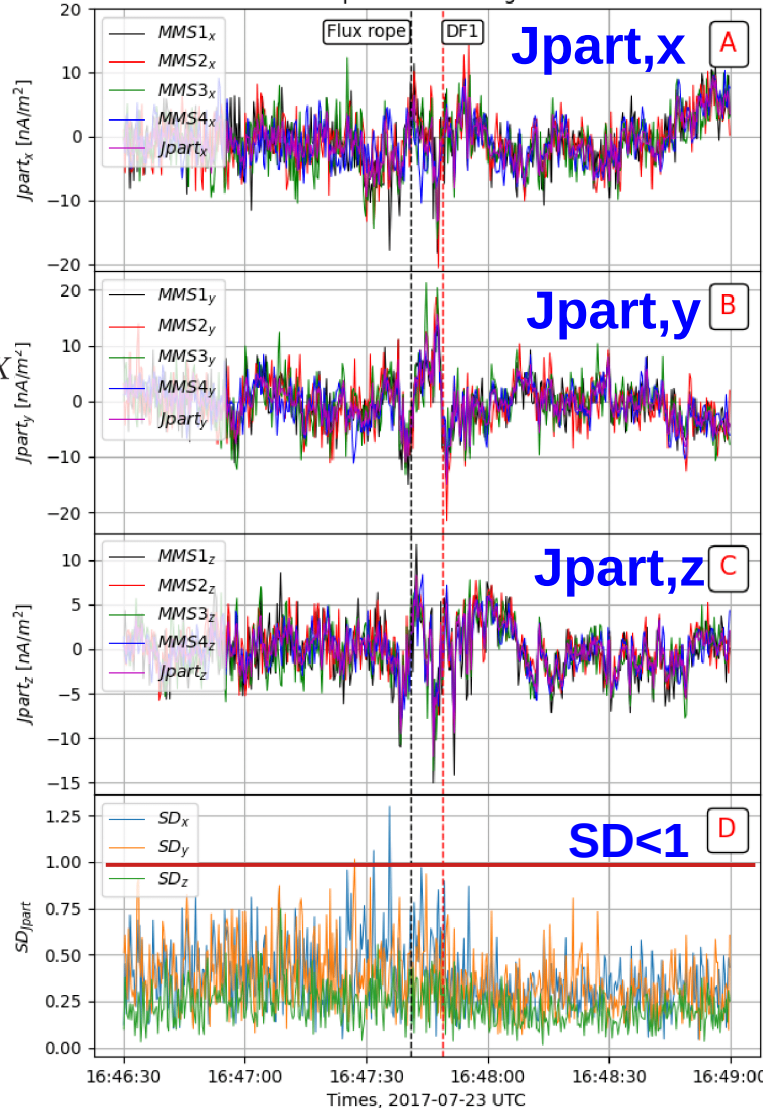


Compute the standard deviation (SD) **normalized** by its error bar:

$$SD(X)/\Delta X = \sqrt{\sum_{i=1}^4 (X_i - \langle X \rangle)^2 / 4} / \Delta X$$

- $\langle X \rangle$: The four spacecraft average of the X component
- ΔX : Respective estimated error bar
- $\Delta E'$ ~1.7 mV/m
- ΔJ ~6.8 nA/m²

Larger **normalized SD** for E' field than for Jpart suggest that non homogeneity comes from the E field.



Case study summary

[Alqeeq et al., PoP, MMS special issue, 2022]



- Good agreement between current densities calculated from particles and curl B. Ions are decoupled at DF mostly due to Hall field but also possibly due to electron pressure gradient assuming a non-zero curl of this term.
- **In the frame of the satellite**, the energy is dissipated ($\mathbf{J} \cdot \mathbf{E} > 0$, dissipation or load region) ahead of the DF but transferred from the plasma to the field behind the front ($\mathbf{J} \cdot \mathbf{E} < 0$, dynamo or generator region).
- **In the fluid frame**, the energy is transferred from the plasma to the fields ($\mathbf{J} \cdot \mathbf{E}' < 0$, dynamo) as also found in a previous MMS single DF event [Yao et al., 2017].
- The energy conversion is **not homogeneous** at the electron scale (scale of the tetrahedron) mostly due to the E field fluctuations which are likely related to LHD waves [e.g., M. Hosner et al. 2022].

